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Quantitative Structure of an Acetate Dye-Molecule Analogue at the TiO₂-Acetic Acid Interface

Hadeel Hussain¹†, Xavier Torrelles³, Gregory Cabailh⁴, Parasmani Rajput²^, Robert Lindsay⁵, Oier Bikondoa⁶, Marcus Tillotson¹, Ricardo Grau-Crespo⁸, Jörg Zegenhagen²‡, Geoff Thornton¹*

- ^{1.} London Centre for Nanotechnology and Department of Chemistry, University College London, 20 Gordon Street, London WC1H OAJ, UK
- ^{2.} ESRF, The European Synchrotron. 71, Avenue des Martyrs, CS40220, F-38043 Grenoble cedex 09, France
- ^{3.} Institut de Ciència de Materials de Barcelona (CSIC), Campus UAB, 08193 Bellaterra, Spain
- ^{4.} Sorbonne Universités, UPMC Univ Paris 06, CNRS-UMR 7588, Institut des NanoSciences de Paris, F-75005, Paris, France
- ^{5.} Corrosion and Protection Centre, School of Materials, The University of Manchester, Sackville Street, Manchester M13 9PL, UK
- ^{6.} Department of Physics, University of Warwick, Gibbet Hill Road, Coventry, CV4 7AL, UK
- ^{7.} XMaS, the UK-CRG Beamline. ESRF, The European Synchrotron. 71, Avenue des Martyrs, CS40220, F-38043 Grenoble cedex 09, France
- 8. Department of Chemistry, University of Reading, Reading RG6 6AD, UK

Present Addresses

- † Corrosion and Protection Centre, School of Materials, The University of Manchester, Sackville Street, Manchester, M13 9PL, UK
- ^ Atomic & Molecular Physics Division, Bhabha Atomic Research Centre, Trombay, Mumbai-400085, India
- ‡ Diamond Light Source Ltd., Diamond House, Harwell Science and Innovation Campus, Didcot, Oxfordshire, OX11 0DE, UK

Corresponding Author

*E-mail: g.thornton@ucl.ac.uk. Tel: +44 (0)20 7679 7979.

Abstract

The positions of atoms in and around acetate molecules at the rutile $TiO_2(110)$ interface with 0.1 M acetic acid have been determined with a precision of ± 0.05 Å. Acetate is used as a surrogate for the carboxylate groups typically employed to anchor monocarboxylate dye molecules to TiO_2 in dye-sensitised solar cells (DSSC). Structural analysis reveals small domains of ordered (2 x 1) acetate molecules, with substrate atoms closer to their bulk terminated positions compared to the clean UHV surface. Acetate is found in a bidentate bridge position, binding through both oxygen atoms to two five-fold titanium atoms such that the molecular plane is along the [001] azimuth. Density functional theory calculations provide adsorption geometries in excellent agreement with experiment. The availability of these structural data will improve the accuracy of charge transport models for DSSC.

Introduction

The interaction of carboxylic acids with TiO₂ is important in a number of applications. For instance, dyes in dye-sensitized solar cells (DSSC) are typically anchored to TiO₂ via one or more carboxylate groups.^{1,2} In this work we use acetate as a surrogate for a dye molecule with a single functionality. However, structural studies have thus far focused on exposure of TiO₂ to carboxylic acids at/near ultra high vacuum (UHV) conditions with a large proportion focused on formic acid. Recently, we used UHV scanning tunneling microscopy (STM) to study the adsorption sites of carboxylates formed at aqueous interfaces, finding that the same sites are occupied as those identified in UHV adsorption.³ Here we move even closer to measurements in a technologically relevant environment, examining in a quantitative fashion the *in situ* structure of the TiO₂(110) interface formed upon immersion in 0.1 M acetic acid using surface X-ray diffraction (SXRD). This concentration is chosen to match that used in an infra-red spectroscopy study of the TiO₂ acetic acid interface, the results of which point to bidentate bonding of acetate³.

It is well established that exposing $TiO_2(110)$ to CH_3COOH vapor at room temperature results in an ordered (2 x 1) overlayer at saturation coverage, i.e. 0.5 monolayers (ML).⁴⁻⁶ The overlayer consists of acetate molecules ([CH₃COO]) formed via deprotonation, which bind to the surface through both oxygen atoms to two adjacent 5-fold surface titanium atoms, so that the molecular plane is aligned with the [001] azimuth. The cleaved H⁺ is thought to adsorb on neighboring bridging oxygen atoms.⁴⁻⁶ A ball and stick model of the adsorption geometry is shown in Figure 1. Density functional theory (DFT) calculations have confirmed that the bidentate adsorption mode is thermodynamically preferred over monodentate adsorption, and that the p(2 x 1) pattern is more stable than the c(2 x 2) pattern with the same concentration.⁷ Our SXRD results from the $CH_3COOH_{(aq)}$ interface with $TiO_2(110)$ show that carboxylate forms small overlayer domains of (2 x 1) symmetry.

Experimental Details

SXRD measurements were carried out using the six-circle diffractometer end station on ID32 at the European Synchrotron Radiation Facility. The $TiO_2(110)$ (Pi-Kem) sample was prepared in UHV using the Sample Characterization Laboratory at ID32 (base pressure $\sim 1 \text{ x} \cdot 10^{-10} \text{ mbar}$). This involved cycles of Ar^+ sputtering at room temperature and 1000 K annealing until a sharp (1 x 1) low energy electron diffraction pattern was obtained and row-resolution imaging in STM was achieved. Auger electron spectroscopy was used to confirm a lack of contamination.

After preparation the sample was transferred under UHV to a small portable ion-pumped UHV chamber ("baby chamber") with a base pressure in the 10⁻⁹ mbar range. This chamber incorporates a dome shaped beryllium window to allow unrestricted transmission of the incident and reflected X-ray beams. Subsequently, this baby chamber was mounted on the ID32 six-circle diffractometer for SXRD measurements, with the sample surface in the horizontal plane. Measurements were performed with the sample at room temperature using a monochromatic focused beam with energy of 17.7 keV ($\lambda = 0.7$ Å), defined by slits to a size of (H x V: 200 µm x 20 µm). The angle of incidence of the X-ray beam with respect to the surface was kept constant at 0.3° for all measurements with 2 x 2 mm² slits in front of a point detector. The experimental data were measured using rocking scans in which scattered intensity is measured whilst the sample is rotated about its normal. These intensities were then integrated and geometrical correction factors¹¹ were applied to evaluate the structure factors which, when represented versus perpendicular momentum transfer, are known as crystal truncation rods (CTRs). The TiO2 (110) surface unit cell was described by lattice vectors $(\mathbf{a_1}, \mathbf{a_2}, \mathbf{a_3})$ parallel to the $[1\bar{1}0]$, [001] and [110] directions, respectively, where $\mathbf{a_1} =$ $a_3 = a\sqrt{2}$ and $a_2 = c$ (a = 4.593 Å and c = 2.958 Å).

A large data set, comprising 1284 non-equivalent reflections, was recorded for the UHV-prepared surface. Immediately afterwards, the baby chamber was vented with N_2 (99.998 % purity) within a glove bag where the sample was transferred to a mylar thin film cell 12, which was then mounted on the ID32 six-circle diffractometer. Next, a 0.1 M aqueous solution of acetic acid was injected into the mylar thin film cell, and a further 1479 non-equivalent reflections were measured. For these SXRD measurements, the mylar thin film cell was in thin-layer-geometry *i.e.* only a thin layer of $CH_3COOH_{(I)}$ (< 1 μ m) was in contact with the sample surface. 12 It should be noted that fractional order rods (FORs) were also investigated but no measurable intensity was found.

Determination of the surface structure involved the use of a least squares fitting procedure implemented within a modified version of ROD, ¹³ in which simulated CTRs are generated. By optimizing the geometry of the structure, a best-fit model between theory and experiment can be found as measured by optimized $\chi^{2 \text{ 14}}$ and R-factor. ¹⁵ A χ^2 value close to 1 and an R-value close to 0.10 is considered an excellent fit between the experimentally observed and the theoretically calculated structure factors. We note that given their low X-ray scattering cross-section, H atoms were ignored in the fitting procedure.

Trial computer simulations for acetate adsorbed on $TiO_2(110)$ (at the vacuum interface) were performed using the density functional theory (DFT) code VASP.¹⁶ Geometries and total energies were obtained using the generalised gradient approximation (GGA) with the exchange-correlation functional of Perdew, Burke and Ernzerhof (PBE).¹⁷ A Hubbard-type correction to the Ti 3d orbitals was applied, with a value of $U_{\rm eff}$ = 3 eV.¹⁸ No van der Waals corrections were applied as their contribution to the strong chemisorption of carboxylic acids at this surface can be expected to be small.⁷ The $TiO_2(110)$ surface was simulated using a periodic stack of quasi-two dimensional slabs, where each slab is separated from the neighbor by a vacuum gap of 15 Å. We used slabs with four O-Ti-O tri-layers, with the bottom two layers fixed to equilibrium bulk positions. A discussion about the adequacy of this model to simulate the adsorption of molecules at this surface can be seen in Ref [¹⁹].

Results and Discussion

SXRD analysis of UHV-prepared TiO₂(110)(1 x 1) considers a surface slab consisting of 2.5 unit cells (6 TiO atomic layers) in the [110] direction that involved a total of 49 parameters: 34 atomic displacements, 12 Debye-Waller (DW) factors, a scale factor, a roughness parameter and a surface fraction parameter. The best-fit model produced a χ^2 and R-value of 1.5 and 0.14, respectively. These values represent a high degree of agreement between experiment and simulated patterns. From the comparison of the atomic displacements from this study and that of a previous SXRD study²⁰ on TiO₂(110), listed in Table S1, it is clear that the two models are essentially in quantitative agreement with each other. This evidences the high reproducibility of the preparation procedure as described in the experimental section. All DW factors adopted reasonable values for both Ti and O atoms where the highest value corresponded to the topmost surface layer ($B_{Ti} = 0.5 \pm 0.1$, $B_O = 0.9 \pm 0.1$) progressively decreasing with depth to their bulk value ($B_{Ti} = 0.3 \pm 0.1$, $B_O = 0.2 \pm 0.1$). The β roughness parameter²¹ adopted a value ($\beta = 0.07$) consistent with the very flat surface seen in STM images recorded prior to SXRD measurements. The surface fraction parameter confirmed that the entire surface adopted the geometry described by the best-fit model.

The optimized TiO₂(110)(2 x 1)-[CH₃COO] structure obtained following analysis of the SXRD data produced a χ^2 of 1.05 and an R-value of 0.15. The total number of parameters used are the same as that for the UHV-prepared surface with an additional 9 parameters associated with the acetate molecule; 6 parameters to determine orientation and site position, 2 DW factors (O and C atoms) and an occupancy parameter. Excellent agreement between the best-fit model and experiment is achieved, as can be seen from the comparison of the experimental and best fit simulated data in Figure 2. This displays experimental CTRs (black error bars), with their best-fit theoretical simulations for the surface after exposure to 0.1 M CH₃COOH_(aq) (red line). It also shows the effect of removing the acetate molecule from the model and reoptimizing the structural and non-structural parameters (*i.e.* DW factors) (blue line). This increases the χ^2 and R-factor to values of 1.56 and 0.19, respectively. It is clearly evident that the presence of the acetate molecule significantly improves the overall goodness-of-fit. The adsorbed acetate was found to be in a bidentate bridge location *i.e.* binding

through both oxygen atoms to two five-fold titanium atoms such that its molecular plane is along the [001] azimuth (Figure 1), in excellent agreement with the literature for acetate and other simple carboxylates bound to this substrate at the vacuum interface. The resulting bond distances are: $Ti(2) - O(11) = 2.13 \pm 0.03$ Å, $Ti(2^*) - O(12) = 2.06 \pm 0.03$ Å, which compare well with bond distances from the formate/ $TiO_2(110)$ system.

In contrast to Refs. 23 and 24, here we allow the adsorbed carboxylate to be asymmetric in order to compare with DFT calculations. The results of the latter indicate an asymmetric carboxylate, which appears to arise from the orientational ordering of the hydroxyl groups. In the experiment, the hydroxyls will likely be disordered and this asymmetry will be averaged, and indeed the nominal asymmetry is almost within the error bars (see Table 1). The expected orientational disorder of the hydroxyl groups was confirmed by DFT calculations in a 2 x 2 supercell, where the results indicate that all different orientational configurations are within 20 meV in energy.

The apparent volume of the molecule has also been investigated²⁵ by simulating a uniform expansion/compression of the molecule. The best fit is consistent with a small molecular expansion (2.9 \pm 0.6 %), however this is considered to be insignificant because of the lack of sensitivity of the χ^2 and R-factor to this change in volume (Table S3 provides optimized positions of atoms in the acetate moiety).

Table 1 lists the internal bond distances and angles of acetate/acetic acid emerging from this study and previous literature values. There is excellent agreement with both previous experimental^{26,27} and theoretical²⁸ work regarding the intra-molecular bond distances. Discrepancies in bond angles are due to previous experimental work being on molecular acetic acid rather than acetate.

When comparing displacements of substrate atoms between the UHV-prepared TiO₂(110)(1 x 1) and the TiO₂(110)(2 x 1)-[CH₃COO] surfaces it is clear that adsorption of acetate reduces the clean surface relaxations (see Tables S1 and S2 for atomic displacements and atomic coordinates, respectively). This almost certainly arises from an increase in substrate surface atom coordination number, and has been seen before for carboxylate/TiO₂(110) at the vacuum interface²³ and other metal oxide surfaces exposed to liquid water.²⁹ This is also reflected in the DW parameters, which return to close to their bulk value. Good agreement of atomic displacements can also be seen between SXRD and DFT-PBE calculations from this study (see supporting information for details). As regards the β roughness parameter, this now takes a value of 0.2, increasing from 0.07. This most likely arises from the small size of randomly distributed domains (~ a few nm) of both ordered acetate moieties and domains absent of molecules, as evidenced in STM images.³ This would also explain the absence of fractional order rods.

In summary, the interface between 0.1 M acetic acid and $TiO_2(110)$ comprises small domains of (2 x 1) ordered arrays of acetate moieties. The positions of the atoms within the molecule as well as the substrate determined by SXRD are in excellent agreement with PBE+U

calculations. It is possible that steric hindrance will prevent the same ordering for monocarboxylic dye molecules in DSSC, for which we are using acetate as a surrogate. However, it seems likely that the same local bonding geometry will be adopted. On this basis, these data should allow more reliable calculations of charge transport at dye molecule-TiO₂(110) interfaces in DSSC.

Supporting Information

Density functional theory calculation results. A ball and stick model of the TiO₂(110)(2 x 1)-[CH₃COO]⁻ surface. DFT-PBE calculated lateral view structure of TiO₂(110)(2 x 1)-[CH₃COO]⁻. Atomic displacements for the as-prepared TiO₂ (110)(1 x 1) and TiO₂(110)(2 x 1)-[CH₃COO]⁻ substrate atoms from SXRD and DFT-PBE. Atomic coordinates of the ideal bulk-terminated coordinate of each atom with the corresponding shift from this position after the fitting procedure. Experimental optimized positions of the atoms in the acetate molecule.

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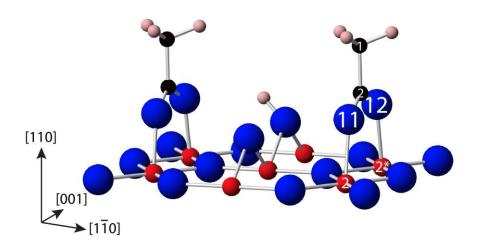


Figure 1 - A ball and stick model of the TiO₂(110)(2 x 1)-[CH₃COO]⁻ surface. In the current study, the adsorbed acetate was found to be in a bidentate bridge location with the cleaved H⁺ thought to adsorb on neighboring bridging oxygen atoms. Large blue, small red, small black and small pink spheres are oxygen, titanium, carbon and hydrogen, respectively. The labeling identifies the atom positions in Tables 1, S1, S2 and S3.

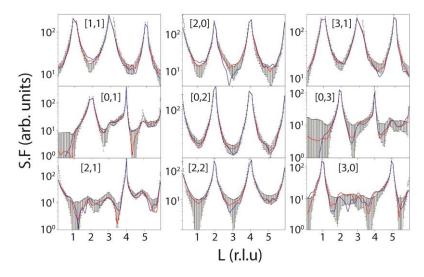


Figure 2 - Experimentally observed (black error bars) and calculated CTRs for the $TiO_2(110)(2 \text{ x 1})$ -[CH₃COO] surface model (red line) and after removing the acetate molecule (blue line). It is clear that the addition of the acetate molecule to the model improves the overall goodness-of-fit.

Table 1 - Comparison of molecular acetic acid (CH₃COOH) bond distances and angles^{26–28} with that of acetate ([CH₃COO]) resulting from SXRD and DFT-PBE in this work. Atom labels are given in Figure 1 and S1.

	Bond Distance (Å) / Angle (°)				
Atoms	SXRD	DFT-PBE	[26]	[27]	[28]
	[CH ₃ COO]	$[CH_3COO]^{-}$	[CH ₃ COOH]	[CH ₃ COOH]	[CH ₃ COOH]
C(1) - C(2)	1.54 ±	1.51	1.49	1.52	1.52
	0.03				
C(2) - O(11)	$1.31 \pm$	1.28	1.31	1.36	1.38
	0.04				
C(2) - O(12)	$1.32 \pm$	1.28	1.23	1.21	1.23
	0.04				
Ti(2) - O(11)	$2.13 \pm$	2.11	-	-	-
	0.03				
Ti(2*) - O(12)	$2.06 \pm$	2.08	-	-	-
	0.03				
Ð C(1)C(2)O(11)	117.7 ±	117.7	-	110.6	111.2
	0.9				
Ð C(1)C(2)O(12)	116.9 ±	117.0	-	126.6	125.9
	0.8				